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Clustering of chiral particles in flows with broken parity invariance

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Motion of an 'isotropic helicoid'

Equations for velocity \mathbf{v} and angular velocity $\boldsymbol{\omega}$ for small isotropic helicoid:

Happel & Brenner, Low Reynolds number hydrodynamics (1963)

$$\dot{\mathbf{v}} = \frac{1}{\tau_p} \left[\mathbf{u}(\mathbf{r}, t) - \mathbf{v} + \frac{2a}{9} C_0 (\boldsymbol{\Omega}(\mathbf{r}, t) - \boldsymbol{\omega}) \right]$$

$$\dot{\boldsymbol{\omega}} = \frac{1}{\tau_p} \left[\frac{10}{3} (\boldsymbol{\Omega}(\mathbf{r}, t) - \boldsymbol{\omega}) + \frac{5}{9a} C_0 (\mathbf{u}(\mathbf{r}, t) - \mathbf{v}) \right]$$

Stokes' law

translation – rotation coupling (scalar)

\mathbf{u} Fluid velocity

$\boldsymbol{\Omega}$ Half fluid vorticity

τ_p Particle relaxation time

$a = \sqrt{5I_0/(2m)}$ Particle 'size' (defined by mass m and moment of inertia I_0)

C_0 Helicoidality

Equations break spatial reflection symmetry ($\boldsymbol{\Omega}$ and $\boldsymbol{\omega}$ pseudovectors)



Dimensionless parameters

Stokes number $St \equiv \frac{\tau_p}{\tau_\eta}$ Size $\bar{a} \equiv \frac{a}{\eta}$ Helicoidality C_0

with τ_η and η smallest time- and length scales of flow.

Constraint on C_0 :

Using $\mathbf{v}_\pm = \mathbf{v} + B\boldsymbol{\omega}$, $\mathbf{u}_\pm = \mathbf{u} + B\boldsymbol{\Omega}$ with $B \equiv \frac{a(21 \pm \sqrt{441 + 40C_0^2})}{10C_0}$

The equations of motion becomes

$$\dot{\mathbf{v}}_+ = \frac{1}{\tau_p} \frac{39 + \sqrt{441 + 40C_0^2}}{18} (\mathbf{u}_+ - \mathbf{v}_+)$$

$$\dot{\mathbf{v}}_- = \frac{1}{\tau_p} \frac{20(27 - C_0^2)}{9(39 + \sqrt{441 + 40C_0^2})} (\mathbf{u}_- - \mathbf{v}_-)$$

Solution blows up unless $-\sqrt{27} < C_0 < \sqrt{27}$

St and \bar{a} constrained by particle density higher than that of the fluid and geometrical size must be smaller than η .

Example of an isotropic helicoid

Recipe from Lord Kelvin:

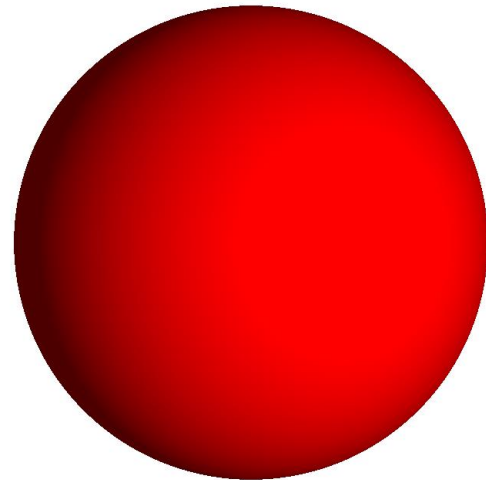
“An isotropic helicoid can be made by attaching projecting vanes to the surface of a globe in proper positions; for instance cutting at 45° each, at the middles of the twelve quadrants of any three great circles dividing the globe into eight quadrantal triangles.”

Kelvin, Phil. Mag. **42** (1871)

Example of an isotropic helicoid

Recipe from Lord Kelvin (1884)

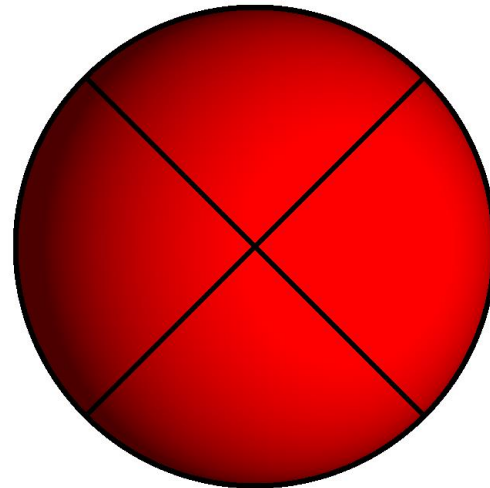
Start with a sphere



Example of an isotropic helicoid

Recipe from Lord Kelvin (1884)

- ✓ Start with a sphere
- Draw 3 great circles

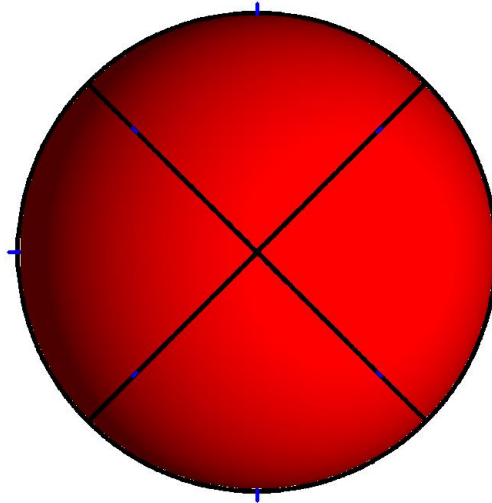


Example of an isotropic helicoid

Recipe from Lord Kelvin (1884)

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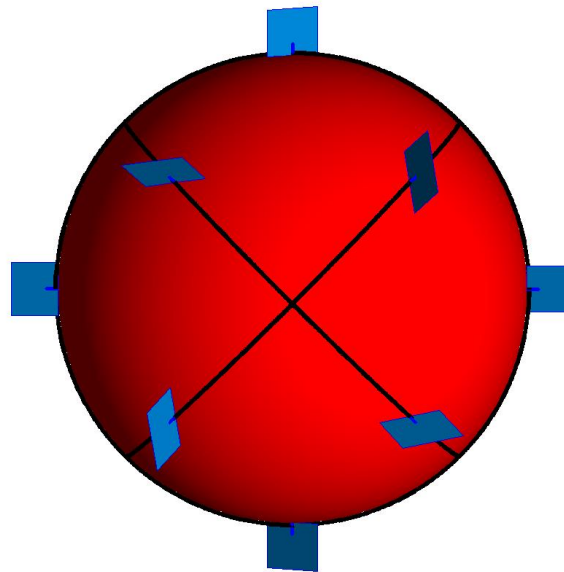
Identify 12 vane positions at midpoints of quarter-arcs



Example of an isotropic helicoid

Recipe from Lord Kelvin (1884)

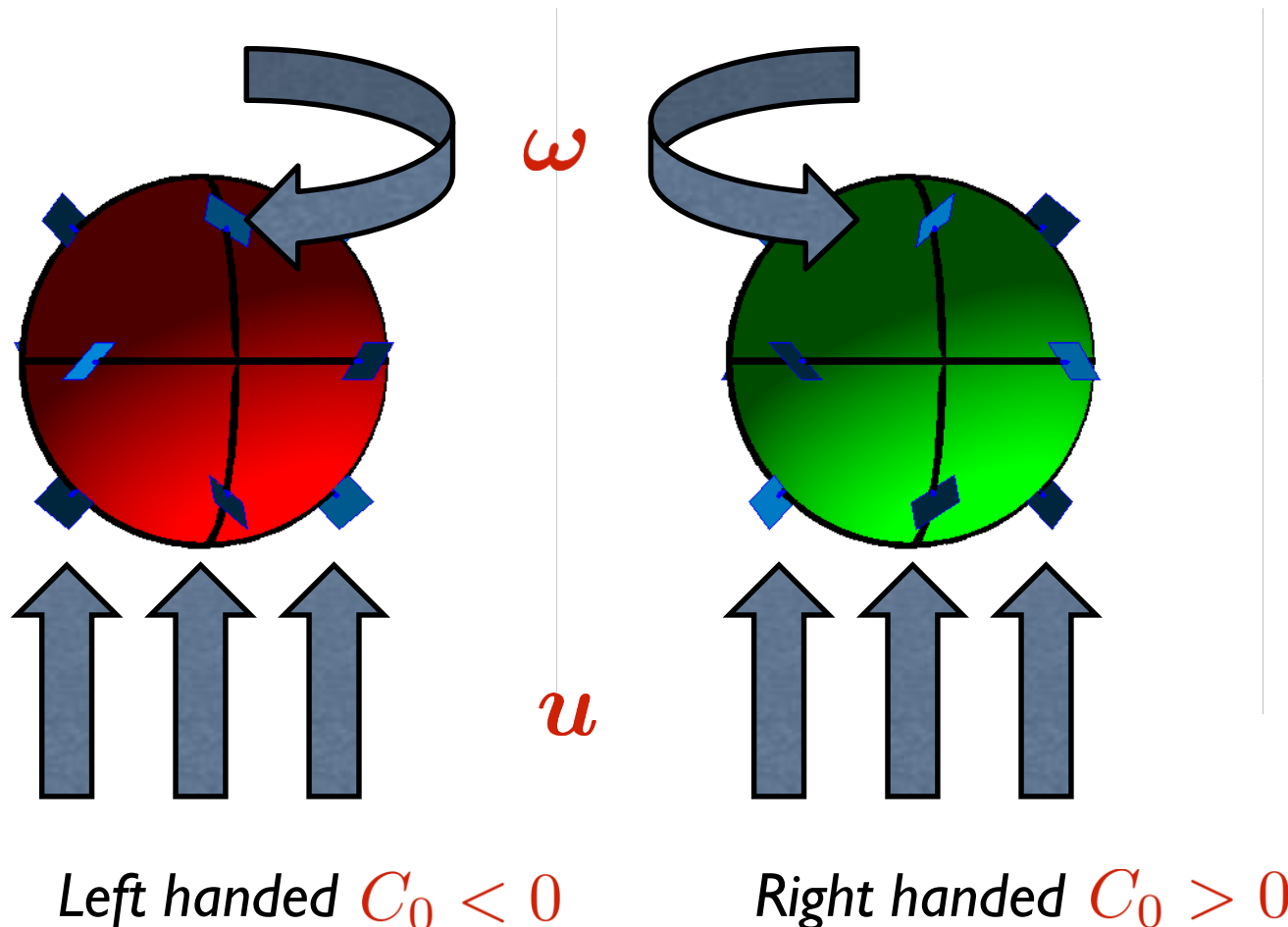
- ✓ Start with a sphere
 - ✓ Draw 3 great circles
 - ✓ Identify 12 vane positions at midpoints of quarter-arcs
- Put a vane on each vane position (45° to arc line)



Chirality

In a constant flow u , the isotropic helicoid starts spinning around the flow direction with angular velocity ω .

The spinning direction depends on the chirality of the vanes.



Clustering at small St

Expand compressibility of particle-velocity field $\nabla \cdot \mathbf{v}$ to first order in St

$$\nabla \cdot \mathbf{v} = -\frac{27}{27 - C_0^2} \tau_p \left[\text{Tr}(\nabla \mathbf{u}^T \nabla \mathbf{u}^T) - \frac{1}{15} a C_0 \text{Tr}(\nabla \mathbf{u}^T \nabla \Omega^T) \right]$$

Centrifuge effect with modified amplitude

Maxey, J. Fluid Mech. **174** (1987)

Term due to parity breaking of system

Reflection-invariant systems have $\langle \text{Tr}(\nabla \mathbf{u}^T \nabla \Omega^T) \rangle = 0$

Isotropic helicoids violate that relation $\langle \text{Tr}(\nabla \mathbf{u}^T \nabla \Omega^T) \rangle \propto \tau_p C_0$

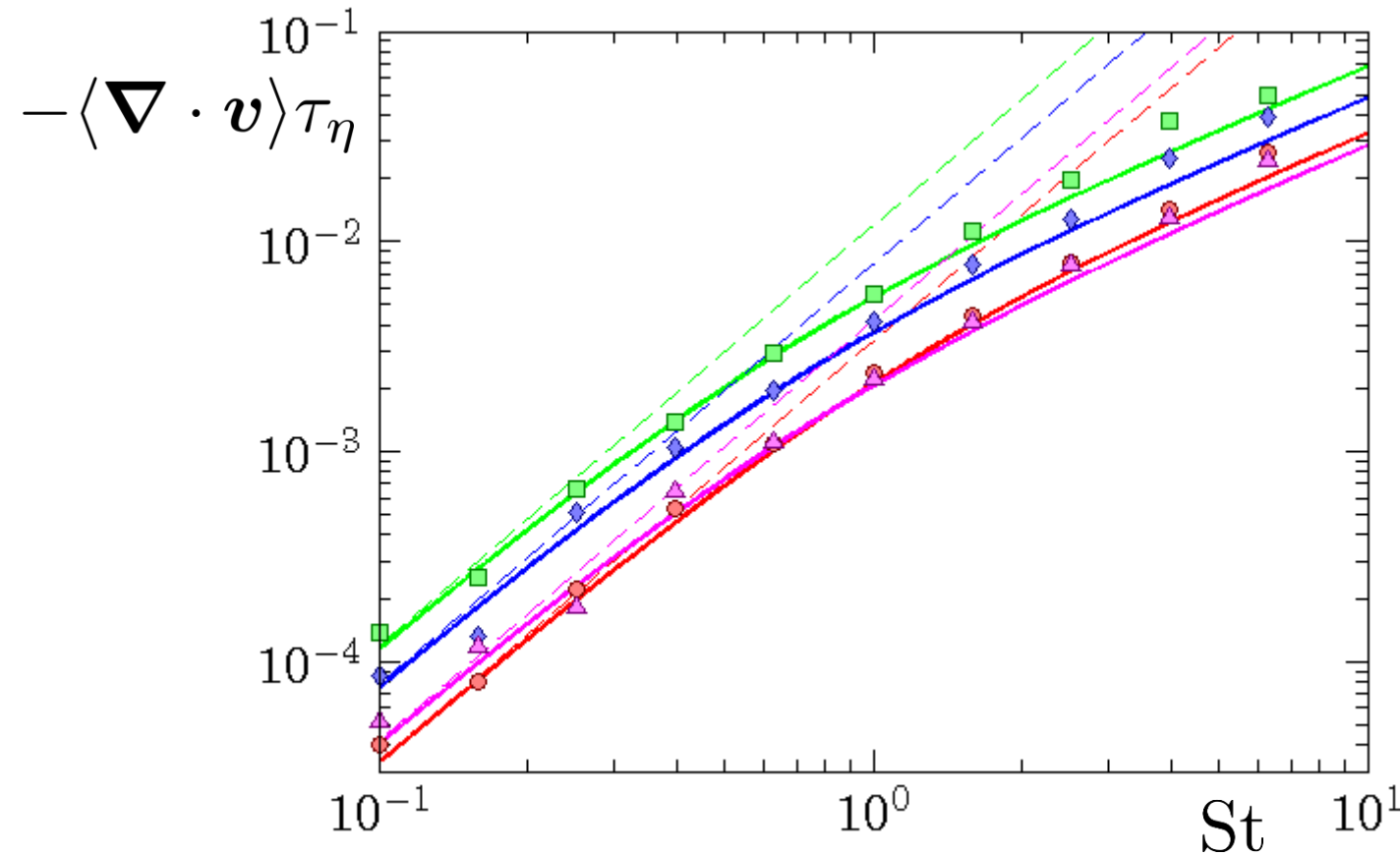
Same for parity-breaking flows $\langle \text{Tr}(\nabla \mathbf{u}^T \nabla \Omega^T) \rangle \propto \tau_p K$

Helicity parameter K

$K > 0$ Right-handed structures ($\mathbf{u} \cdot \Omega > 0$) more common

$K < 0$ Left-handed structures ($\mathbf{u} \cdot \Omega < 0$) more common

Clustering at small St in random flow



$$Ku \equiv \frac{u_0 \tau_\eta}{\eta} = 0.1$$

$$\bar{a} = 1$$

Small- Ku theory
 Gustavsson & Mehlig EPL **96** (2011)

Small- St limit

- Spherical particle ($C_0 = 0$) in neutral flow
- Right-handed particle ($C_0 = 3$) in left-handed flow
- ◆ Right-handed particle ($C_0 = 3$) in neutral flow
- ▲ Right-handed particle ($C_0 = 3$) in left-handed flow

Clustering at small St in random flow

$$\langle \nabla \cdot \mathbf{v} \rangle \tau_\eta = - \frac{27 Ku^4 St^2}{13(27 - C_0^2)(27(10 + 13 St + 3 St^2) - 10C_0^2)^3} \left[\frac{6656}{5\pi} \bar{a}^2 C_0^2 K^2 (27 - C_0^2) (27(10 + 39 St + 15 St^2) - 10C_0^2) \right. \\ - 192 \bar{a} C_0 K \sqrt{2/\pi} (1300 C_0^4 - 27 C_0^2 (2600 + 5070 St + 2457 St^2 + 324 St^3) + 729 (1300 + 5070 St + 4654 St^2 + 1845 St^3 + 351 St^4 + 27 St^5)) \\ + 4550 \bar{a}^2 C_0^6 + 852930 (10 + 3 St)^3 (1 + 3 St + St^2) - 45 C_0^4 (21 \bar{a}^2 (260 + 507 St + 195 St^2 + 18 St^3) - 20 (1300 + 27 St^3)) \\ \left. + 243 C_0^2 (10 + 3 St) (21 \bar{a}^2 (65 + 234 St + 247 St^2 + 87 St^3 + 9 St^4) - 10 (2600 + 4290 St + 1677 St^2 + 261 St^3 + 27 St^4)) \right]$$