On the edge of an inverse cascade

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Forward and inverse cascades

There are some systems ...



- Fast rotating flows ($Ro \equiv U/\Omega \ell \ll 1$)
- Flows in the presence of a magnetic field ($M \equiv U/B_0 \ll 1$)
- Confined flows (thin geometries) ($\Gamma \equiv h/\ell_f \ll 1$)
- Helical MHD flows ($h_M \equiv$ helicity injection/energy injection $\cdot k_f$)

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for which the inverse cascade depends on a parameter

$$\mu = Ro, \Gamma, M, h_M, \dots$$

Motivation

Fluxes in: Thin layers/Rotating/Stratified/Magnetic fields ...



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A turbulence to turbulence transition ...



- the system transitions from one turbulent state (inverse cascading) to an other (forward cascading) varying a parameter μ. (μ is not Re)
- the transition occurs in the presence of turbulent noise
- these transitions are not only observed as dimensional (ie 2D to 3D), but weak to strong, HD to MHD, ...
- these transitions are not only observed for the energy cascade but also for other invariants (magnetic helicity, square vector potential, wave action, ...)

Most models have a varying inverse cascade of energy due to a dimensional transition from a 3D to a 2D

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Are there computationally more tractable models that show a transition from forward to inverse cascade?

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Thus they require 3D high resolution numerical simulations.

Are there computationally more tractable models that show a transition from forward to inverse cascade? Transitioning from 2D-HD to 2D-MHD





2D-HD vs 2D-MHD

Equations:

$$\partial_t \omega + \mathbf{u} \cdot \nabla \omega = + \left[\nu_n^+ \nabla^{2n} \omega + \nu_n^- \nabla^{-2n} \omega\right] + F_\omega$$

where

$$\omega = \hat{\mathbf{e}}_n \cdot \nabla \times \mathbf{u}$$

Nonlinearity conserves

$$E = \frac{1}{2} \int \mathbf{u}^2 dv, \qquad \Omega = \frac{1}{2} \int \omega^2 dv$$

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2D-HD vs 2D-MHD

Equations:

$$\partial_t \omega + \mathbf{u} \cdot \nabla \omega = \mathbf{b} \cdot \nabla j + [\nu_n^+ \nabla^{2n} \omega + \nu_n^+ \nabla^{-2n} \omega] + F_{\omega}$$

$$\partial_t a + \mathbf{u} \cdot \nabla a = + [\eta_n^+ \nabla^{2n} a + \eta_n^- \nabla^{-2n} a] + F_a$$

where

$$\boldsymbol{\omega} = \hat{\mathbf{e}}_n \cdot \nabla \times \mathbf{u}, \quad \mathbf{b} = \nabla \times (\hat{\mathbf{e}}_n a), \quad j = \hat{\mathbf{e}}_n \cdot \nabla \times \mathbf{b}$$

Nonlinearity conserves

$$E = \frac{1}{2} \int \mathbf{u}^2 + \mathbf{b}^2 dv, \qquad A = \frac{1}{2} \int a^2 dv$$

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2D-HD vs 2D-MHD



 $\mathbf{F}_b = 0 \Rightarrow \mathbf{b} = \mathsf{No}$ dynamo theorem for 2D flows!

$ \mathbf{F}_u > 0, \mathbf{F}_b = 0$	$ \mathbf{F}_u > 0, \mathbf{F}_b > 0$	$\mathbf{F}_u = 0, \mathbf{F}_b > 0$
Inverse cascade of E	?	Forward cascade of E
Forward cascade of Ω	?	not conserved
Forward cascade of A	?	Inverse cascade of A

What is the fate of the forward/inverse cascade as we vary F_u, F_b ?

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Set-up of Numerical Experiments



2D square periodic box of side $2\pi L$ No mean magnetic field $\langle {\bf b} \rangle = {\bf 0}$

$$F_{\omega}(x,y) = f_u k_f^{+1} \sin(k_f x) \sin(k_f x)$$

$$F_a(x,y) = f_b k_f^{-1} \cos(k_f x) \cos(k_f x)$$

hypodissipation $\nu_n^- \nabla^{-2n}$ & hyperdissipation $\nu_n^+ \nabla^{+2n}$ with n=2. Control Parameters / Non-dimensional Numbers $\mu_f \equiv \frac{f_b}{f_u} \qquad k_f L \qquad Re_f^- = \frac{f_u^{1/2} k_f^{1/2+2n}}{\nu_n^-} \qquad Re_f^+ = \frac{f_u^{1/2} k_f^{1/2-2n}}{\nu_n^+}$ $P_M^- \equiv \nu_n^- / \eta_n^- = 1, \qquad P_M^+ \equiv \nu_n^+ / \eta_n^+ = 1$

Quantifying the cascades

Inverse and Forward cascades of energy:

$$\begin{split} \epsilon_{\scriptscriptstyle E}^- &\equiv \nu_n^- \langle (\nabla^{-n} \mathbf{u})^2 + (\nabla^{-n} \mathbf{b})^2 \rangle, \qquad \epsilon_{\scriptscriptstyle E}^+ \equiv \nu_n^+ \langle (\nabla^{+n} \mathbf{u})^2 + (\nabla^{+n} \mathbf{b})^2 \rangle \\ \epsilon_{\scriptscriptstyle E} &\equiv \epsilon_{\scriptscriptstyle E}^- + \epsilon_{\scriptscriptstyle E}^+ \qquad \qquad 0 \leq \frac{\epsilon_{\scriptscriptstyle E}^-}{\epsilon_{\scriptscriptstyle E}} \leq 1, \end{split}$$

Inverse and Forward cascades of square vector potential:

$$\begin{split} \epsilon_A^- &\equiv \nu_n^- \langle (\nabla^{-n} a)^2 \rangle, \qquad \epsilon_A^+ \equiv \nu_n^+ \langle (\nabla^{+n} a)^2 \rangle \\ \epsilon_A^- &\equiv \epsilon_E^- + \epsilon_E^+ \qquad \qquad 0 \leq \frac{\epsilon_A^-}{\epsilon_A} \leq 1, \end{split}$$

Varying μ_f for different box-size and fixed Re_n^+ .



A Critical transition



critical behavior:

$$\epsilon_E^- \propto (\mu_{c_E} - \mu)^{\gamma_E}$$
 and $\epsilon_A^- \propto (\mu - \mu_{c_A})^{\gamma_A}$

a best fit leads to:

 $\mu_{c_E} \simeq 0.22 \dots, \ \gamma_E \simeq 0.82 \quad and \quad \mu_{c_A} \simeq 0.25 \dots, \ \gamma_A \simeq 0.27$ Э

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Critical point dependence on Re_n^+

Varying μ_f for different values of Re_n^+ and fixed box-size.



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$$\mu_c = \mu_c(Re_n^+)$$

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Critical point dependence on Re_n^+ (rescaling)

Varying μ_f for different values of Re_n^+ and fixed Re_n^- .



• $\mu_c \propto (Re_n^+)^{-1/2n}$

Critical point dependence on Re_n^+ (rescaling)

Varying μ_f for different values of Re_n^+ and fixed Re_n^- .



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$$\mu_c \propto (Re_n^+)^{-1/2n}$$

Magnetic tension determines the transition:

$$\mu_b \equiv \frac{b^2 k_f}{f_u} \propto \mu_f^2 \left(\frac{k_d^+}{k_f}\right)^2 \propto \mu_f^2 [Re^+]^{1/n}$$

Large scale spectra



Small scale spectra

Varying μ_f for $Re_n^+ \gg 1$.



For small μ

- magnetic energy at the smallest scales is $b_\ell^2 \propto \mu^2 \ell_d^{-2}$ (passive advectio)
- kinetic energy at the smallest scales is $u_{\ell}^2 \propto \epsilon_{\Omega}^{2/3} \ell_d^2$ (enstrophy cascade)

Nonlinearity starts when

$$\mu \geq \mu_{\scriptscriptstyle NL} \propto \ell_d^2 \propto R e^{-1/n}$$

Variable forward and backward fluxes



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Instantaneous and time averaged fluxes

Strong fluctuations of the energy fluxes



$$\mu = 0.21 \dots \lesssim \mu_c$$



Kinetic energy

Magnetic energy

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Thank you for your attention!

 $\mu \ll \mu_{\scriptscriptstyle NL}$



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Current density

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 $\mu_{\scriptscriptstyle NL} \lesssim \mu \ll \mu_c$



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Current density

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 $\mu_{\scriptscriptstyle NL} \ll \mu \ll \mu_c$



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 $\mu_{NL} \ll \mu \lesssim \mu_c$



Vorticity

Current density

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$$\mu = 0.21 \dots \lesssim \mu_c$$



Kinetic energy

Magnetic energy

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$$\mu = 0.21 \dots \lesssim \mu_c$$



Kinetic energy

Vector Potential

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$$\mu = 0.26 \cdots \gtrsim \mu_c$$



Kinetic energy

Magnetic energy

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$$\mu = 0.26 \cdots \gtrsim \mu_c$$



Kinetic energy

Vector Potential

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Small scale dissipations



For small μ

- magnetic energy at the smallest scales is $b_\ell^2 \propto \mu^2 \ell_d^{-2}$ (passive advectio)
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A cartoon summary



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Break down of the enstrophy conservation



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Digging in in the transition



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